

2. Designate the corners of a square, l on a side, in clockwise order, A, B, C, D . Put charges $2q$ at A and $-3q$ at B . Determine the value of the line integral of \mathbf{E} , from point C to point D . (No actual integration needed!) What is the numerical answer if $q = 10^{-9} C$ and $l = 5\text{cm}$? [2.5]

Assignment 6 (Sols.)

1. Excess charge, $Q = 5 \times 10^8 \times 1.6 \times 10^{-19} \text{C} = 8 \times 10^{-11} \text{C}$.

~~Charge density~~, $\sigma = \frac{Q}{l} = \frac{8 \times 10^{-11} \text{C}}{0.04 \text{m}} = 2 \times 10^{-9} \text{C/m}^2$.

(a) Electric field in the rest frame,

$$E = \frac{\sigma}{2\pi\epsilon_0 r} \approx \frac{2 \times 10^{-9} \text{C/m}^2}{(0.0001)\text{m}} \times 18 \times 10^9 \frac{\text{Nm}^2}{\text{C}^2}$$

$$\approx 72 \times 10^4 \text{V/m}. \quad (\because \frac{1}{4\pi\epsilon_0} \approx 9 \times 10^9 \frac{\text{Nm}^2}{\text{C}^2}, \quad \therefore \frac{1}{2\pi\epsilon_0} \approx 18 \times 10^9 \frac{\text{Nm}^2}{\text{C}^2}).$$

Direction : radial.

(b) In the moving frame,

$$E'_r = \gamma E_r \quad \text{where,} \quad \gamma = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}} = \frac{1}{\sqrt{1 - (0.9)^2}} \approx 2.29.$$

$$= 2.29 \times 72 \times 10^4 \text{V/m}$$

$$\approx 1.65 \times 10^6 \text{V/m}$$

direction: radial.

PHY102 : Quiz 1

1. A spherical charge distribution has a density ρ that is constant from $r = 0$ out to $r = R$ and is zero beyond. What is the electric field for all values of r , both less than and greater than R ? [2.5]

2. Since the charged particle feels the force due to the electric field only in the y -direction, therefore

$$\therefore f_x = \frac{d p_x}{dt} = 0 \Rightarrow p_x \text{ is conserved.}$$

However, p is relativistic momentum and is given as $\vec{p} = \gamma m_0 \vec{u}$ where m_0 : rest mass

$$\gamma = \frac{1}{\sqrt{1 - u^2/c^2}}$$

Since p_x is conserved,

$$\therefore (p_x)_{\text{before entering field}} = (p_x)_{\text{after entering field.}}$$

$$(p_x)_b = \gamma_b m_0 (u_x)_b = \frac{m_0 (u_x)_b}{\sqrt{1 - \frac{(u_x)_b^2}{c^2}}}.$$

$$(p_x)_a = \gamma_a m_0 (u_x)_a = \frac{m_0 (u_x)_a}{\sqrt{1 - \frac{(u_x)_a^2 + (u_y)_a^2}{c^2}}}$$

Since, the particle entering the region of the electric field has a y -component of velocity as well.

$$\therefore \frac{m_0 (u_x)_b}{\sqrt{1 - \frac{(u_x)_b^2}{c^2}}} = \frac{m_0 (u_x)_a}{\sqrt{1 - \frac{(u_x)_a^2 + (u_y)_a^2}{c^2}}}.$$

$$\Rightarrow (u_x)_b^2 \left(1 - \frac{(u_x)_a^2 + (u_y)_a^2}{c^2} \right) = (u_x)_a^2 \left(1 - \frac{(u_x)_b^2}{c^2} \right)$$

PHY102 : Quiz 1

1. A spherical charge distribution has a density ρ that is constant from $r = 0$ out to $r = R$ and is zero beyond. What is the electric field for all values of r , both less than and greater than R ? [2.5]

$$\rightarrow \frac{(\mathbf{U}_n)_b^2 - (\mathbf{U}_n)_a^2}{c^2} - \frac{(\mathbf{U}_n)_b (\mathbf{U}_n)_a}{c^2} = (\mathbf{U}_n)_a^2 - \frac{(\mathbf{U}_n)_a (\mathbf{U}_n)_b}{c^2}$$

$$\Rightarrow (\mathbf{U}_n)_a^2 = (\mathbf{U}_n)_b^2 \left[1 - \frac{(\mathbf{U}_n)_a}{c} \right].$$

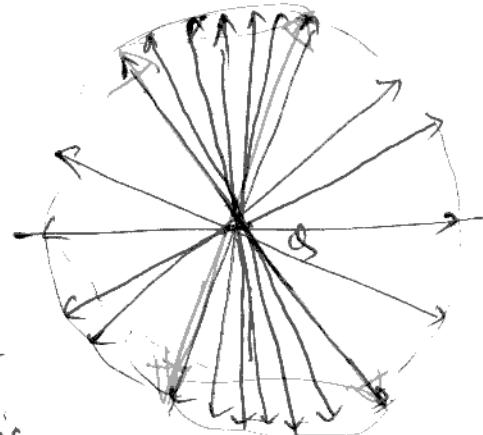
$$\therefore (\mathbf{U}_n)_a = (\mathbf{U}_n)_b \sqrt{1 - \frac{(\mathbf{U}_n)_a}{c}}.$$

Since, $(\mathbf{U}_n)_a \neq 0$, $\therefore (\mathbf{U}_n)_a < (\mathbf{U}_n)_b$

$\therefore n$ component of velocity decreases.

3. The electric field due to a moving charge looks as in the figure - symmetric about a plane \perp to direction of charge but not spherically symmetric.

Field stronger at right angles than in direction of motion. This helps understand why half of total flux Φ could be contained between 2 conical surfaces as shown in red lines.

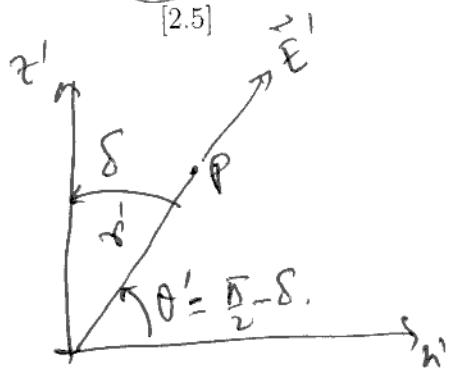


2. Designate the corners of a square, l on a side, in clockwise order, A, B, C, D . Put charges $2q$ at A and $-3q$ at B . Determine the value of the line integral of \mathbf{E} , from point C to point D . (No actual integration needed!) What is the numerical answer if $q = 10^{-9} \text{ C}$ and $l = 5\text{cm}$? [2.5]

$$\text{Now, } E' = \frac{1}{4\pi\epsilon_0} \frac{Q}{r'^2} \frac{1-\beta^2}{(1-\beta^2 \sin^2 \theta')^{3/2}}$$

If flux through 2 conical surfaces is half the total flux from Q , then by symmetry,

flux from Q between $\theta' = 0$ & $\frac{\pi}{2} - \delta$ = flux from Q between $\theta' = \frac{\pi}{2} - \delta$ to $\frac{\pi}{2}$.



$$\text{Now, flux} = E' \cdot 2\pi r'^2 \sin \theta' d\theta'.$$

$$\therefore \int_0^{\pi/2 - \delta} \frac{1}{4\pi\epsilon_0} \frac{Q}{r'^2} \frac{(1-\beta^2) \cdot 2\pi r'^2 \sin \theta' d\theta'}{(1-\beta^2 \sin^2 \theta')^{3/2}},$$

$$= \int_{\pi/2 - \delta}^{\pi} \frac{1}{4\pi\epsilon_0} \frac{Q}{r'^2} \frac{(1-\beta^2) 2\pi r'^2 \sin \theta' d\theta'}{(1-\beta^2 \sin^2 \theta')^{3/2}}.$$

$$\Rightarrow \int_0^{\pi/2 - \delta} \frac{\sin \theta' d\theta'}{(1-\beta^2 \sin^2 \theta')^{3/2}} = \left\{ \begin{array}{l} \int_{\pi/2 - \delta}^{\pi} \\ \end{array} \right. \frac{\sin \theta' d\theta'}{(1-\beta^2 \sin^2 \theta')^{3/2}}.$$

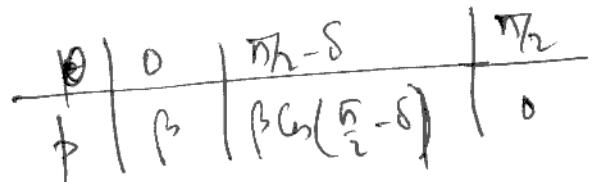
$$\Rightarrow \int_0^{\pi/2 - \delta} \frac{\sin \theta' d\theta'}{(1-\beta^2 + \beta^2 \cos^2 \theta')^{3/2}} = \left\{ \begin{array}{l} \int_{\pi/2 - \delta}^{\pi} \\ \end{array} \right. \frac{\sin \theta' d\theta'}{(1-\beta^2 + \beta^2 \cos^2 \theta')^{3/2}}$$

$$(\because \sin^2 \theta' = 1 - \cos^2 \theta').$$

PHY102 : Quiz 1

1. A spherical charge distribution has a density ρ that is constant from $r = 0$ out to $r = R$ and is zero beyond. What is the electric field for all values of r , both less than and greater than R ? [2.5]

Let, $\Phi = \beta \cos\theta' \quad \therefore d\Phi = -\beta \sin\theta' d\theta'$



$$\int_{\beta}^{\beta \cos(\frac{\pi}{2} - \delta)} \frac{-d\Phi/\beta}{(1 - \beta^2 + \Phi^2)^{3/2}} = \int_{\beta \cos(\frac{\pi}{2} - \delta)}^0 \frac{-d\Phi/\beta}{(1 - \beta^2 + \Phi^2)^{3/2}}$$

$$\Rightarrow \int_{\beta \cos(\frac{\pi}{2} - \delta)}^0 \frac{d\Phi}{(1 - \beta^2 + \Phi^2)^{3/2}} = \int_0^{\beta \cos(\frac{\pi}{2} - \delta)} \frac{d\Phi}{(-\beta^2 + \Phi^2)^{3/2}}$$

Since, $\int \frac{dx}{(a^2 + x^2)^{3/2}} = \frac{x}{a^2(a^2 + x^2)^{1/2}}$ (prob. S. 11)

\therefore taking $a^2 = 1 - \beta^2$,

$$\left[\frac{1}{(1 - \beta^2)(1 - \beta^2 + \Phi^2)^{1/2}} \right]_{\beta \cos(\frac{\pi}{2} - \delta)}^0 = \left[\frac{1}{(1 - \beta^2)(1 - \beta^2 + \Phi^2)^{1/2}} \right]_{0}^{\beta \cos(\frac{\pi}{2} - \delta)}$$

$$\Rightarrow \beta - \frac{\beta \cos(\frac{\pi}{2} - \delta)}{\left[1 - \beta^2 \sin^2(\frac{\pi}{2} - \delta) \right]^{1/2}} = \frac{\beta \cos(\frac{\pi}{2} - \delta)}{\left[1 - \beta^2 \sin^2(\frac{\pi}{2} - \delta) \right]^{1/2}} - 0.$$

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$$\therefore f = \frac{2\beta \sin \delta}{[1 - \beta^2 \cos \delta]^n}$$

$$\Rightarrow 1 - \beta^2 \cos \delta = 4 \sin^2 \delta.$$

$$\Rightarrow 1 - \beta^2 + \beta^2 \sin^2 \delta = 4 \sin^2 \delta.$$

$$\Rightarrow (4 - \beta^2) \sin^2 \delta = 1 - \beta^2$$

$$\therefore \sin \delta = \pm \sqrt{\frac{1 - \beta^2}{4 - \beta^2}}$$

$$= \pm \frac{1}{r} \left(\frac{1}{4 - \beta^2} \right)^{1/2} \quad \left(\because r^2 = \frac{1}{1 - \beta^2} \right)$$

For $r \gg 1$, δ is small $\therefore \sin \delta \approx \delta$.

Also, $\beta \approx 1$ for large r .

$$\therefore \delta \approx \pm \frac{1}{r} \frac{1}{\sqrt{3}}$$

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4. In the rest frame of the protons, the electrostatic force between the 2 protons is just $\frac{e^2}{4\pi\epsilon_0 r^2}$.

$$\therefore f_{\text{rest}} = \frac{e^2}{4\pi\epsilon_0 r^2}.$$

Now, getting back to the lab frame, which is moving with velocity βc ,

$$f = \frac{1}{\gamma} f_{\text{rest}} = \frac{1}{\gamma} \frac{e^2}{4\pi\epsilon_0 r^2}.$$

However, at the instantaneous position of one of the protons, the electric field strength caused by the other is $\frac{\gamma e}{4\pi\epsilon_0 r^2}$.

$$\therefore \text{Discrepancy} = \frac{1}{\gamma} \frac{e^2}{4\pi\epsilon_0 r^2} - \left(\frac{\gamma e}{4\pi\epsilon_0 r^2} \right) \cdot e$$

$$= \frac{e^2}{4\pi\epsilon_0 r^2} \left(1 - \frac{1}{\gamma} \right).$$

$$= \frac{e^2}{4\pi\epsilon_0 r^2} \cancel{\gamma} \left(\frac{1}{\gamma} - 1 \right).$$

$$= -\frac{e^2}{4\pi\epsilon_0 r^2} \gamma \beta^2 = -\frac{e^2}{4\pi\epsilon_0 r^2} \cdot \gamma \frac{v}{c} \beta.$$

$$= -ev \underbrace{\left(\frac{\gamma e}{4\pi\epsilon_0 r^2} \cdot \frac{\beta}{c} \right)}_{B \text{ (magnetic field)}}$$

$$\therefore B = \frac{\gamma e}{4\pi\epsilon_0 r^2} \frac{\beta}{c} = \frac{\gamma}{c} E$$

Note: B is (β/c) times the electric field & not β as in question.

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5. In frame F , $I_k = \gamma_k \beta_k c$.

In the frame F' , $\beta'_k = \frac{\beta_k + \beta}{1 + \beta_k \beta}$ since F' moves w/ to line with velocity $(-\beta c)$.

Now, $\gamma'_k = \frac{\gamma_k \gamma'_k}{\gamma_k}$ (recall the argument from what I did in the ~~last~~ class.
 $F \rightarrow \text{rest frame} \rightarrow F'$).

$$\gamma_k = \frac{1}{\sqrt{1 - \beta_k^2}} ; \quad \gamma'_k = \frac{1}{\sqrt{1 - \beta'_k^2}}, \quad \gamma = \frac{1}{\sqrt{1 - \beta^2}}$$

$$\therefore \gamma'_k = \frac{1}{\sqrt{1 - \left(\frac{\beta_k + \beta}{1 + \beta_k \beta}\right)^2}} = \gamma \gamma_k \left(1 + \beta \beta_k\right).$$

$$\begin{aligned} \therefore I'_k &= \gamma'_k \beta'_k c = \gamma_k \frac{\gamma \gamma_k (1 + \beta \beta_k)}{\gamma_k} \left(\frac{\beta_k + \beta}{1 + \beta \beta_k}\right) c \\ &= \gamma_k \gamma (\beta_k + \beta) c = \gamma (\gamma_k \beta_k c + \gamma_k \beta c) \\ &= \gamma (I_k + \beta c \gamma_k). \end{aligned}$$

$$\lambda' = \gamma_k \cdot \frac{\gamma \gamma_k (1 + \beta \beta_k)}{\gamma_k} = \gamma \left(\gamma_k + \beta \frac{I_k}{c} \right)$$

$$\text{Total, } \lambda = \sum_n \lambda_n \quad \& \quad I = \sum_n I_n$$

$$\therefore \lambda' = \gamma \left(\lambda + \beta \frac{I}{c} \right) ; \quad I' = \gamma \left(I + \beta c \lambda \right).$$